

## A Sensitive Method for the Measurement of Three-Bond $C'$ , $H^\alpha$ $J$ Couplings in Uniformly $^{13}C$ - and $^{15}N$ -Enriched Proteins

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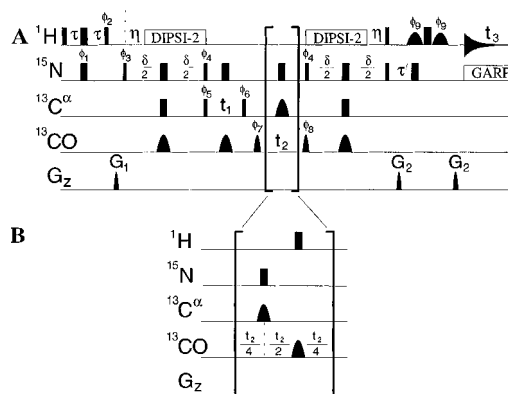
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Vicinal spin–spin coupling constants can provide direct information about local conformation. A variety of NMR methods have been developed to determine the magnitude of homonuclear and heteronuclear couplings in isotopically labeled proteins.<sup>1</sup> Due to the ambiguity inherent to Karplus-type relations,<sup>2</sup> accurate torsion angles can only be derived from a combination of several coupling constants. For the polypeptide backbone angle  $\phi$ , the measurement of  $^3J(C,H)$  coupling constants<sup>3–5</sup> supplements the structural information available from  $^3J(H^N,H^\alpha)$ . Recently, we have introduced an E.COSY-type<sup>6</sup> experiment for the quantitative determination of  $^3J(C'_{i-1},H^\alpha_i)$ .<sup>4</sup> A drawback of this (H)NCAHA pulse sequence is that  $\alpha$ -protons are detected during acquisition while the use of  $H_2O$  as the solvent is mandatory. Hence, the evaluation of coupling constants is impossible for residues with  $^1H^\alpha$  resonances in the vicinity of the intense water signal. In the HCAN-[C']E.COSY described by Wang and Bax,<sup>5</sup> this difficulty is circumvented by a magnetization transfer pathway suitable for samples dissolved in  $D_2O$ . Thus, accurate values for  $^3J(C'_{i-1},H^\alpha_i)$  can, in principle be obtained for all residues in a protein. However, as a consequence of the long periods of  $^{13}C^\alpha$  transverse magnetization, this pulse sequence is expected to be relatively inefficient for proteins with fast  $R_{2,C^\alpha}$  relaxation rates. In this paper, we propose a sensitive E.COSY-type experiment,  $H^\alpha$ -coupled H(N)CA,CO, in which the  $^3J(C'_{i-1},H^\alpha_i)$  displacement occurs in an indirectly detected  $^{13}C'$  dimension, thereby avoiding difficulties caused by the solvent resonance.

The new method is based on the COHNCA,<sup>7</sup> a dimensionality reduced sequence which simultaneously correlates amide protons and nitrogens with carbonyl carbons of the preceding and  $\alpha$ -protons of the same and the preceding residue. In the pulse scheme depicted in Figure 1,  $^{13}C^\alpha$  and  $^{13}C'$  chemical shifts are sampled in two independent dimensions. Since no proton decoupling is applied during  $t_1$  and  $t_2$ , the large  $^1J(C^\alpha,H^\alpha)$  coupling leads to a splitting of the signals along the  $F_1$  ( $^{13}C^\alpha$ ) dimension. Vicinal and geminal  $^{13}C'$ ,  $^1H^\alpha$  couplings evolving in the subsequent  $t_2$  ( $^{13}C'$ ) domain can be extracted from the



**Figure 1.** Pulse scheme of the  $H^\alpha$ -coupled H(N)CA,CO experiment without (version A) and with (version B) chemical shift scaling in the  $^{13}C'$  dimension. Narrow and wide pulses denote 90 and 180° flip angles, respectively. The proton carrier frequency is placed on the water resonance and is shifted to the center of the amide region for the DIPSI-2<sup>12</sup> decoupling sequences. The carbon carrier position is 58 ppm until completion of the  $t_1$  evolution and changed to 176 ppm following the pulse with the phase  $\phi_6$ . The shaped  $^{13}C^\alpha$  pulse is a G3 Gaussian cascade<sup>13</sup> with a duration of 0.5 ms. Carbonyl 90 and 180° pulses have a duration of 124  $\mu s$  and an amplitude profile of the center lobe of a sinc function. GARP  $^{15}N$  decoupling<sup>14</sup> during acquisition is performed using a 0.75 kHz radiofrequency field. Solvent suppression is achieved with the WATERGATE method,<sup>15</sup> applied during the  $^1H, ^{15}N$  reverse INEPT step.<sup>16</sup> The width of the 90° Gaussian-shaped pulses for selective excitation of the water resonance is 2.5 ms. Gradient durations and strengths are the following  $G_1$ , 1 ms and 10 G  $cm^{-1}$ ;  $G_2$ , 0.8 ms and 35 G  $cm^{-1}$ . Phase cycling is as follows:  $\phi_1 = x, -x$ ;  $\phi_2 = y, -y$ ;  $\phi_3 = 2(x), 2(-x)$ ;  $\phi_4 = y$ ;  $\phi_5 = 4(x), 4(-x)$ ;  $\phi_6 = x + 25^\circ$ ;  $\phi_7 = 8(x), 8(-x)$ ;  $\phi_8 = x + 47^\circ$ ;  $\phi_9 = -x$ ; receiver  $x, 2(-x), x, -x, 2(x), 2(-x), 2(x), -x, x, 2(x), x$ . Nonlabeled pulses are applied along the  $x$ -axis. The phases  $\phi_6$  and  $\phi_8$  are adjusted to compensate for zero-order Bloch–Siegert phase errors as indicated. Quadrature in the  $t_1$  and  $t_2$  domains is obtained by changing the phases  $\phi_5$  and  $\phi_7$ , respectively, in the States–TPPI manner.<sup>17</sup> Delay durations are  $\tau = 2.3$  ms,  $\delta = 28.6$  ms,  $\eta = 5.4$  ms, and  $\tau' = 2.7$  ms. Acquisition times were 14.1 ms in  $t_1$  and 106.5 ms in  $t_3$ . In  $t_2$ , 82 ( $t_{2,max} = 75.4$  ms) and 52 ( $t_{2,max} = 95.7$  ms) complex increments were recorded for A and B, respectively. Accumulation of 16 scans per FID resulted in measuring times of 72 h (A) and 41 h (B).

displacement of the two components of the intra- and interresidual correlations, respectively (Figure 2A). It might be argued that the assignment of the role of the passive spin to  $\alpha$ -protons leads to an underestimation of the  $^{13}C', ^1H^\alpha$  coupling constants due to different relaxation rates of antiphase terms compared to those of in-phase terms.<sup>8</sup> This effect is, however, minimized as the time between  $^{13}C^\alpha$  and  $^{13}C'$  evolution periods, in which  $^1H^\alpha$  spin flips can occur, is essentially zero owing to the synchronous buildup of  $^{15}N$  antiphase magnetization with respect to both carbon species.

Sufficient resolution in the  $^{13}C'$  domain requires a large number of  $t_2$  increments, leading to a relatively long measuring time. Although extensive aliasing would afford relief, we prefer to employ chemical shift scaling<sup>9</sup> to avoid the accidental introduction of signal overlap from different spectral regions. In version B of the  $H^\alpha$ -coupled H(N)CA,CO pulse sequence, the evolution of two-bond and three-bond  $^{13}C', ^1H^\alpha$  couplings during  $t_2$  is scaled up by a factor of 2 with respect to  $^{13}C'$  chemical shifts (Figure 2B). As a result, the number of increments can be substantially decreased, while retaining the

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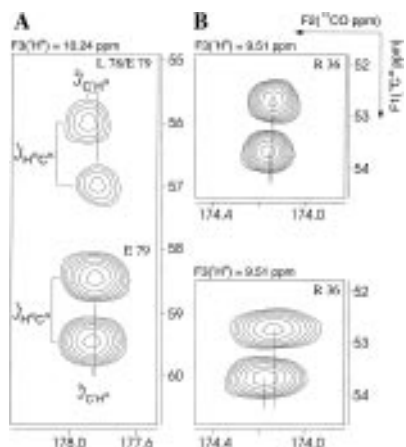
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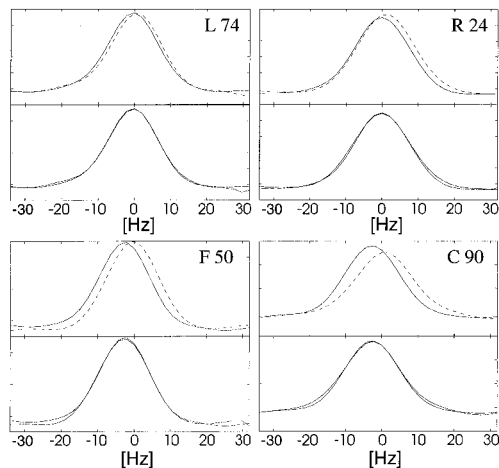
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**Figure 2.** (A) Expansion of a  $F_1, F_2$  plane of the  $H^\alpha$ -coupled  $H(N)$ -CA,CO spectrum (version B) at the  $F_3$  ( $^1H^N$ ) chemical shift of residue Glu 79 from *D. vulgaris* flavodoxin. The horizontal displacements of the two doublet components are due to  $^3J(C'_{i-1}, H^\alpha_i)$  and  $^2J(C'_{i-1}, H^\alpha_{i-1})$  couplings for intraresidual and sequential crosspeaks, respectively. (B) Effect of chemical shift scaling. The intraresidual crosspeaks of Arg 36 in the spectra recorded with versions A (top) and B (bottom) of the  $H(N)$ CA,CO pulse sequence are shown. Since identical ppm scales are used, the line width and the E.COSY splitting appear doubled in the second case.

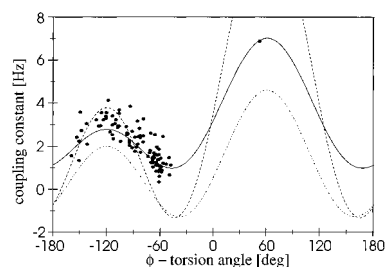


**Figure 3.** Traces from the  $H^\alpha$ -coupled  $H(N)$ CA,CO spectrum (version B) of flavodoxin, obtained by summation over the spectral points in  $F_1$  and  $F_3$ . The upper panels represent the original traces through the low-field (solid line) and high-field (dashed line) multiplet halves. In the lower panels, a superposition of the fitted traces is shown. A continuous frequency shift is achieved by inverse Fourier transformation, followed by multiplication of the  $t_2$  time domain data with a complex phase factor.<sup>10</sup> The  $^3J(C'_{i-1}, H^\alpha_i)$  values resulting from the fitting procedure are  $0.83 \pm 0.12$  Hz (Leu 74),  $1.22 \pm 0.1$  Hz (Arg 24),  $2.41 \pm 0.16$  Hz (Phe 50), and  $3.39 \pm 0.14$  Hz (Cys 90).

spectral width and the accuracy of the frequency determination of the individual crosspeak components.

The two versions of the novel pulse sequence were applied to a 1.4 mM uniformly  $^{13}C/^{15}N$ -enriched sample of oxidized *Desulfovibrio vulgaris* flavodoxin (MW 16.3 kDa). Both experiments were performed at a Bruker DMX 600 spectrometer equipped with a pulsed field gradient unit and an actively shielded gradient triple-resonance probe. For the evaluation of  $J$ -coupling constants, a least-squares fit was carried out in the  $t_2$  time domain as described by Schmidt et al.<sup>10</sup> Representative examples are shown in Figure 3.

Three-bond  $^{13}C', ^1H^\alpha$  couplings were determined for 103 residues of flavodoxin. The root mean square pairwise difference for the coupling constants obtained with versions A and



**Figure 4.** Experimental  $^3J(C'_{i-1}, H^\alpha_i)$  coupling constants versus  $\phi$ -angles derived from the crystal structure of oxidized *D. vulgaris* flavodoxin at 1.7 Å resolution.<sup>18</sup> The corresponding Karplus curves are taken from Wang and Bax<sup>5</sup> (---), Bystrov et al.<sup>19</sup> (---), and Solkan and Bystrov<sup>20</sup> (- · - ·).

B was 0.54 Hz. A graphical representation of the dependence on the  $\phi$ -angle together with the available Karplus curves is given in Figure 4. These values are taken from the spectrum recorded with the chemical shift scaled version of the  $H(N)$ -CA,CO sequence, unless overlap resulting from the broader lines in  $F_2$  prevented the evaluation, as was the case for nine residues. The Karplus parameters of Wang and Bax<sup>5</sup> yielded the best agreement between measured and expected coupling constants although for residues with  $\phi$ -angles around  $-120^\circ$  slightly higher values are obtained in the present study. A relatively large scattering of  $^3J$  couplings is observed for some residues with similar  $\phi$ -angles. Apart from experimental errors, this might be due to a varying amplitude of fluctuations leading to  $J$  averaging<sup>11</sup> or due to differences between the backbone conformation in the crystal and the solution structure. The majority of residues exhibiting large deviations of experimental  $^3J(C'_{i-1}, H^\alpha_i)$  coupling constants with respect to the Karplus curve<sup>5</sup> is located outside regions of regular secondary structure elements. Sequential correlations yielding  $^2J(C'_{i-1}, H^\alpha_{i-1})$  values could be detected for all residues unless obscured by stronger, overlapping signals. Since the intensities are considerably smaller compared to those of intraresidual peaks, a lower precision can be expected for the geminal coupling constants.

In conclusion, a method has been introduced which allows the determination of  $^3J(C'_{i-1}, H^\alpha_i)$  for all non-glycine and non-proline residues in isotopically labeled proteins. The high sensitivity of the pulse sequence should enable its application to larger proteins than the one investigated in this work.

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